Rotating-Coil Scanner (Toy Train) and the Translating Fluxmeter for Extracting Pseudo Multipoles in Accelerator Magnets

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## **Superconducting LHC Magnets**

Round aperture Long and nearly straight (2D) Coil dominated















#### **Iron Dominated Magnets**



3D Eddy currents and hysteresis -> strange wipe-out properties

Combined function Curved Fast ramped







1. Governing equation in the air domain

 $\nabla^2 A_z = 0,$ 

2. Chose a suitable coordinate system

$$r^2 \frac{\partial^2 A_z}{\partial r^2} + r \frac{\partial A_z}{\partial r} + \frac{\partial^2 A_z}{\partial \varphi^2} = 0,$$



3. Make a guess, look it up in a book, use the method of separation: That is: find eigenfunctions. Coefficients are not know yet

$$A_z(r,\varphi) = \sum_{n=1}^{\infty} r^n (\mathcal{A}_n \sin n\varphi + \mathcal{B}_n \cos n\varphi).$$





## 4. Derive the radial field component

$$B_r(r,\varphi) = \frac{1}{r} \frac{\partial A_z}{\partial \varphi} = \sum_{n=1}^{\infty} n r^{n-1} (\mathcal{A}_n \cos n\varphi - \mathcal{B}_n \sin n\varphi),$$

5. Measure (or calculate) the field on a reference radius and perform Fourier analysis (develop into the eigenfunctions). Coefficients are known here.

$$B_r(r_0,\varphi) = \sum_{n=1}^{\infty} (B_n(r_0)\sin n\varphi + A_n(r_0)\cos n\varphi),$$





6: Compare the known and unknown coefficients

$$B_r(r,\varphi) = \frac{1}{r} \frac{\partial A_z}{\partial \varphi} = \sum_{n=1}^{\infty} nr^{n-1} (\mathcal{A}_n \cos n\varphi - \mathcal{B}_n \sin n\varphi),$$
$$B_r(r_0,\varphi) = \sum_{n=1}^{\infty} (B_n(r_0) \sin n\varphi + A_n(r_0) \cos n\varphi),$$
$$\mathcal{A}_n = \frac{1}{n r_0^{n-1}} A_n(r_0), \qquad \mathcal{B}_n = \frac{-1}{n r_0^{n-1}} B_n(r_0).$$

7. Put this into the original solution for the entire air domain

$$A_z(r,\varphi) = -\sum_{n=1}^{\infty} \frac{r_0}{n} \left(\frac{r}{r_0}\right)^n (B_n(r_0)\cos n\varphi - A_n(r_0)\sin n\varphi).$$







#### 8: Calculate fields and potentials in the entire air domain

$$A_z(r,\varphi) = -\sum_{n=1}^{\infty} \frac{r_0}{n} \left(\frac{r}{r_0}\right)^n (B_n(r_0)\cos n\varphi - A_n(r_0)\sin n\varphi).$$

$$B_r(r,\varphi) = \sum_{n=1}^{\infty} \left(\frac{r}{r_0}\right)^{n-1} \left(B_n(r_0)\sin n\varphi + A_n(r_0)\cos n\varphi\right)$$
$$B_{\varphi}(r,\varphi) = \sum_{n=1}^{\infty} \left(\frac{r}{r_0}\right)^{n-1} \left(B_n(r_0)\cos n\varphi - A_n(r_0)\sin n\varphi\right)$$

$$B_x(r,\varphi) = \sum_{n=1}^{\infty} \left(\frac{r}{r_0}\right)^{n-1} (B_n(r_0)\sin(n-1)\varphi + A_n(r_0)\cos(n-1)\varphi)$$
$$B_y(r,\varphi) = \sum_{n=1}^{\infty} \left(\frac{r}{r_0}\right)^{n-1} (B_n(r_0)\cos(n-1)\varphi - A_n(r_0)\sin(n-1)\varphi)$$







Take any  $2\pi$  periodic function and develop according to

$$\frac{C_0}{2} + \sum_{n=1}^{\infty} (C_n(r_0)\sin n\varphi + D_n(r_0)\cos n\varphi).$$

	$B_r$	$B_arphi$	$B_x$	$B_y$	$A_z$	$\phi_{ m m}$
$B_n =$	$C_n$	$D_n$	$C_{n-1}$	$D_{n-1}$	$\frac{-nD_n}{r_0}$	$\frac{-n\mu_0 C_n}{r_0}$
$A_n =$	$D_n$	$-C_n$	$D_{n-1}$	$-C_{n-1}$	$\frac{nC_n}{r_0}$	$rac{-n\mu_0 D_n}{r_0}$

We can use fields, potentials, fluxes, or wire-oscillation amplitudes as "raw data".







## **Rotating Coil Measurements**





Tangential coil Radial flux Radial coil Tangential flux



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х









Local transverse harmonics calculated at different reference radii and scaled with the 2D laws

$$b_n(r_1) = \left(\frac{r_1}{r_0}\right)^{n-N} b_n(r_0),$$

wrong



## **Integrated Harmonics**

$$\nabla^2 \phi_{\mathrm{m}}(x, y, z) = \frac{\partial^2 \phi_{\mathrm{m}}(x, y, z)}{\partial x^2} + \frac{\partial^2 \phi_{\mathrm{m}}(x, y, z)}{\partial y^2} + \frac{\partial^2 \phi_{\mathrm{m}}(x, y, z)}{\partial z^2} = 0.$$
$$\overline{\phi}_{\mathrm{m}}(x, y) := \int_{-z_0}^{z_0} \phi_{\mathrm{m}}(x, y, z) \mathrm{d}z.$$

$$\begin{aligned} \frac{\partial^2 \overline{\phi}_{\mathrm{m}}(x,y)}{\partial x^2} + \frac{\partial^2 \overline{\phi}_{\mathrm{m}}(x,y)}{\partial y^2} &= \int_{-z_0}^{z_0} \left( \frac{\partial^2 \phi_{\mathrm{m}}}{\partial x^2} + \frac{\partial^2 \phi_{\mathrm{m}}}{\partial y^2} \right) \mathrm{d}z \\ &= \int_{-z_0}^{z_0} \left( -\frac{\partial^2 \phi_{\mathrm{m}}}{\partial z^2} \right) \mathrm{d}z = -\left. \frac{\partial \phi_{\mathrm{m}}}{\partial z} \right|_{-z_0}^{z_0} \\ &= H_z(-z_0) - H_z(z_0) \stackrel{!}{=} 0. \end{aligned}$$

## The 2D scaling laws hold for the integrated harmonics





Experimental characterization (magnetic compatibility, electrical interference, rotation quality) of a rotating coil transducer for local multipole scanning



Prototype transport system





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$$\phi_{\rm m}(r,\varphi,z) = \left\{ \begin{array}{c} \cos n\varphi \\ \sin n\varphi \end{array} \right\} I_n(pr) \left\{ \begin{array}{c} \cos pz \\ \sin pz \end{array} \right\}$$

$$I_n(pr) = \sum_{k=0}^{\infty} \frac{1}{k! \, \Gamma(k+n+1)} \left(\frac{pr}{2}\right)^{n+2k}$$

$$\phi_{\mathrm{m}} = \sum_{k=0}^{\infty} \sum_{n=1}^{\infty} r^{n+2k} (\mathcal{C}_{n+2k,n}(z) \sin n\varphi + \mathcal{D}_{n+2k,n}(z) \cos n\varphi)$$





$$\begin{aligned} &\frac{1}{r}\frac{\partial}{\partial r} \Big\{ \sum_{k=0}^{\infty} \sum_{n=1}^{\infty} (n+2k)r^{n+2k} \left( \mathcal{C}_{n+2k,n}(z)\sin n\varphi + \mathcal{D}_{n+2k,n}(z)\cos n\varphi \right) \Big\} \\ &\quad - \frac{1}{r^2} \sum_{k=0}^{\infty} \sum_{n=1}^{\infty} n^2 r^{n+2k} \left( \mathcal{C}_{n+2k,n}(z)\sin n\varphi + \mathcal{D}_{n+2k,n}(z)\cos n\varphi \right) \\ &\quad + \sum_{k=0}^{\infty} \sum_{n=1}^{\infty} r^{n+2k} \left( \mathcal{C}_{n+2k,n}^{(2)}(z)\sin n\varphi + \mathcal{D}_{n+2k,n}^{(2)}(z)\cos n\varphi \right) \\ &= \sum_{k=0}^{\infty} \sum_{n=1}^{\infty} (n+2k)^2 r^{n+2k-2} \left( \mathcal{C}_{n+2k,n}(z)\sin n\varphi + \mathcal{D}_{n+2k,n}(z)\cos n\varphi \right) \\ &\quad - \sum_{k=0}^{\infty} \sum_{n=1}^{\infty} n^2 r^{n+2k-2} \left( \mathcal{C}_{n+2k,n}(z)\sin n\varphi + \mathcal{D}_{n+2k,n}(z)\cos n\varphi \right) \\ &\quad + \sum_{k=1}^{\infty} \sum_{n=1}^{\infty} r^{n+2k-2} \left( \mathcal{C}_{n+2k,n}^{(2)}(z)\sin n\varphi + \mathcal{D}_{n+2k,n}^{(2)}(z)\cos n\varphi \right) \\ &\quad = 0, \end{aligned}$$







$$\mathcal{C}_{n+2k,n}(z)\left((n+2k)^2 - n^2\right) + \mathcal{C}_{n+2k-2,n}^{(2)}(z) = 0,$$
  
$$\mathcal{D}_{n+2k,n}(z)\left((n+2k)^2 - n^2\right) + \mathcal{D}_{n+2k-2,n}^{(2)}(z) = 0.$$

$$\mathcal{C}_{n+2k,n}(z) = \frac{1}{\prod_{m=1}^{k} (n^2 - (n+2m)^2)} \mathcal{C}_{n,n}^{(2k)}(z),$$





$$\begin{split} \phi_{\rm m} &= \sum_{n=1}^{\infty} \left\{ \sum_{k=0}^{\infty} \frac{1}{\prod_{m=1}^{k} (n^2 - (n+2m)^2)} \, \mathcal{C}_{n,n}^{(2k)}(z) \right\} r^n \sin n\varphi \\ &+ \sum_{n=1}^{\infty} \left\{ \sum_{k=0}^{\infty} \frac{1}{\prod_{m=1}^{k} (n^2 - (n+2m)^2)} \, \mathcal{D}_{n,n}^{(2k)}(z) \right\} r^n \cos n\varphi \,, \end{split}$$

$$\begin{split} \phi_{\rm m} &= \sum_{n=1}^{\infty} \left\{ \mathcal{C}_{n,n}(z) - \frac{\mathcal{C}_{n,n}^{(2)}(z)}{4(n+1)} r^2 \\ &+ \frac{\mathcal{C}_{n,n}^{(4)}(z)}{32(n+1)(n+2)} r^4 - \frac{\mathcal{C}_{n,n}^{(6)}(z)}{384(n+1)(n+2)(n+3)} r^6 + \dots \right\} r^n \sin n\varphi \\ &+ \sum_{n=1}^{\infty} \left\{ \mathcal{D}_{n,n}(z) - \frac{\mathcal{D}_{n,n}^{(2)}(z)}{4(n+1)} r^2 \\ &+ \frac{\mathcal{D}_{n,n}^{(4)}(z)}{32(n+1)(n+2)} r^4 - \frac{\mathcal{D}_{n,n}^{(6)}(z)}{384(n+1)(n+2)(n+3)} r^6 + \dots \right\} r^n \cos n\varphi \,, \end{split}$$





## The Leading Term is NOT the Measured One







# **Field Components from Pseudo-Multipoles**

$$\begin{split} \phi_{\mathrm{m}}(r,\varphi) &= \sum_{n=1}^{\infty} r^{n} (\widetilde{\mathcal{C}}_{n}(r,z) \sin n\varphi + \widetilde{\mathcal{D}}_{n}(z) \cos n\varphi) \,. \\ B_{r}(r,\varphi,z) &= -\mu_{0} \sum_{n=1}^{\infty} r^{n-1} (\overline{\mathcal{C}}_{n}(r,z) \sin n\varphi + \overline{\mathcal{D}}_{n}(r,z) \cos n\varphi) \,, \\ B_{\varphi}(r,\varphi,z) &= -\mu_{0} \sum_{n=1}^{\infty} n \, r^{n-1} (\widetilde{\mathcal{C}}_{n}(r,z) \cos n\varphi - \widetilde{\mathcal{D}}_{n}(r,z) \sin n\varphi) \,, \\ B_{z}(r,\varphi,z) &= -\mu_{0} \sum_{n=1}^{\infty} r^{n} \left( \frac{\partial \widetilde{\mathcal{C}}_{n}(r,z)}{\partial z} \sin n\varphi + \frac{\partial \widetilde{\mathcal{D}}_{n}(r,z)}{\partial z} \cos n\varphi \right) \,, \end{split}$$

$$\overline{\mathcal{C}}_{n}(r,z) = n \,\mathcal{C}_{n,n}(z) - \frac{(n+2)\mathcal{C}_{n,n}^{(2)}(z)}{4(n+1)}r^{2} + \frac{(n+4)\mathcal{C}_{n,n}^{(4)}(z)}{32(n+1)(n+2)}r^{4} - \dots$$
$$\widetilde{\mathcal{C}}_{n}(r,z) := \mathcal{C}_{n,n}(z) - \frac{\mathcal{C}_{n,n}^{(2)}(z)}{4(n+1)}r^{2} + \frac{\mathcal{C}_{n,n}^{(4)}(z)}{32(n+1)(n+2)}r^{4} - \dots,$$





$$B_n(r_0,z) = -\mu_0 r_0^{n-1} \overline{\mathcal{C}}_n(r_0,z) = -\mu_0 r_0^{n-1} \left( n \,\mathcal{C}_{n,n}(z) - \frac{(n+2)\mathcal{C}_{n,n}^{(2)}(z)}{4(n+1)} r_0^2 + \frac{(n+4)\mathcal{C}_{n,n}^{(4)}(z)}{32(n+1)(n+2)} r_0^4 - \dots \right) \,.$$







$$\mathcal{F}\{\mathcal{C}_{n,n}(z)\} = \frac{-\mathcal{F}\{B_n(r_0,z)\} \ \mathcal{F}\{K_n(r_0,z)\}}{\mu_0 r_0^{n-1} \left(n - \frac{(n+2)(i\omega)^2}{4(n+1)} r_0^2 + \frac{(n+4)(i\omega)^4}{32(n+1)(n+2)} r_0^4 - \dots\right)}$$



















## **Search Coils Must be Saddle-Shaped**







# **Translating Fluxmeter**







TE Technology Department



# The Translating Fluxmeter (Prototype)











### **Longitudinal Profile Measurements**





















## **Solenoid Field Transducer**





$$B_{z}(r,z) = -\mu_{0} \left( \mathcal{C}_{0,0}^{(1)}(z) - \frac{\mathcal{C}_{0,0}^{(3)}(z)}{4}r^{2} + \frac{\mathcal{C}_{0,0}^{(5)}(z)}{64}r^{4} \right)$$
$$B_{r}(r,z) = -\mu_{0} \left( -\frac{\mathcal{C}_{0,0}^{(2)}(z)}{2}r + \frac{\mathcal{C}_{0,0}^{(4)}(z)}{16}r^{3} \right)$$
$$\mathcal{F}\{B_{r}(r_{0},z)\} = -\mu_{0} \mathcal{F}\{\mathcal{C}_{0,0}(z)\} \left( -\frac{(i\omega)^{2}}{2}r_{0} + \frac{(i\omega)^{4}}{16}r_{0}^{3} \right)$$





- ➔ Iso-perimetric (saddle-shaped) search coils
- ➔ Different analogue bucking scheme as radius scaling is no more valid
- ➔ No feed-down correction for axis misaligment
- ➔ Induced voltages in connectors, wiring
- ➔ Accuracy in longitudinal coil positioning
- ➔ Interference with martensitic/conducting elements in the transducer
- ➔ Effect of measurement errors on deconvolution and field reconstruction



